

MEE5114 Advanced Control for Robotics

Lecture 4: Exponential Coordinate of Rigid Body Configuration

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Outline

- Exponential Coordinate of $SO(3)$ *rotation matrix*
- Euler Angles and Euler-Like Parameterizations
- Exponential Coordinate of $SE(3)$ *hom o transformation matrix*
- Instantaneous Velocity of Moving Frames

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Towards Exponential Coordinate of $SO(3)$

- Recall the polar coordinate system of the complex plane: \mathbb{C}
 - Every complex number $z = x + jy = \rho e^{j\phi}$
 - Cartesian coordinate $(x, y) \leftrightarrow$ polar coordinate (ρ, ϕ)
 - For some applications, the polar coordinate is preferred due to its geometric meaning.

$$\Rightarrow \{ R^T R = I, \det(R) = 1 \}$$

- For any rotation matrix $R \in SO(3)$, it turned out $R = e^{[\hat{\omega}]\theta}$
 - $\hat{\omega}$: unit vector representing the axis of rotation
 - θ : the degree of rotation
 - $\hat{\omega}\theta$ is called the **exponential coordinate** for $SO(3)$.

Exponential Coordinate of $SO(3)$

Proposition 1 (Exponential Coordinate $\leftrightarrow SO(3)$).

- For any unit vector $[\hat{\omega}] \in so(3)$ and any $\theta \in \mathbb{R}$,

$$e^{[\hat{\omega}]\theta} \in SO(3) \quad \leftarrow \text{valid rotation matrix}$$

- For any $R \in SO(3)$, there exists $\hat{\omega} \in \mathbb{R}^3$ with $\|\hat{\omega}\| = 1$ and $\theta \in \mathbb{R}$ such that

$$R = e^{[\hat{\omega}]\theta}$$

$$\text{exp: } [\hat{\omega}]\theta \in so(3) \quad \rightarrow \quad R \in SO(3)$$

$$\text{log: } R \in SO(3) \quad \rightarrow \quad [\hat{\omega}]\theta \in so(3)$$

- The vector $\hat{\omega}\theta$ is called the *exponential coordinate* for R
- The exponential coordinates are also called the *canonical coordinates* of the rotation group $SO(3)$

Rotation Matrix as Forward Exponential Map

- Exponential Map: By definition

$$e^{[\omega]\theta} = I + \theta[\omega] + \frac{\theta^2}{2!}[\omega]^2 + \frac{\theta^3}{3!}[\omega]^3 + \dots$$

- **Rodrigues' Formula:** Given any unit vector $[\hat{\omega}] \in so(3)$, we have

$$e^{[\hat{\omega}]\theta} = I + [\hat{\omega}]\sin(\theta) + [\hat{\omega}]^2(1 - \cos(\theta))$$

Fact: if $\|\hat{\omega}\|=1$, then $[\hat{\omega}] = -[\hat{\omega}]^T$, $[\hat{\omega}]^3 = -[\hat{\omega}]$, $[\hat{\omega}]^4 = [\hat{\omega}]^3[\omega] = -[\omega]^2$

$$\begin{aligned} e^{[\hat{\omega}]\theta} &= I + \underbrace{[\hat{\omega}]\theta}_{\sin \theta} + \frac{\theta^2}{2}[\hat{\omega}]^2 + \frac{\theta^3}{3!}(-[\hat{\omega}]) + \frac{\theta^4}{4!}(-[\hat{\omega}]^2) + \dots \\ &= I + \underbrace{\left(\theta - \frac{\theta^3}{3!} + \frac{\theta^5}{5!} + \dots\right)}_{\sin \theta} [\hat{\omega}] + \underbrace{\left(\frac{\theta^2}{2!} - \frac{\theta^4}{4!} + \frac{\theta^6}{6!} - \dots\right)}_{1 - \cos \theta} [\hat{\omega}]^2 \end{aligned}$$

Examples of Forward Exponential Map

- Rotation matrix $R_x(\theta)$ (corresponding to $\hat{x}\theta$)

$$\text{Rot}(\hat{x}, \theta) = e^{[\hat{x}]\theta} = \mathbf{I} + \sin\theta \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & -1 \\ 0 & 1 & 0 \end{bmatrix} + (1 - \cos\theta) \begin{bmatrix} 0 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & -1 \end{bmatrix}$$

$$\hat{x} = \begin{bmatrix} 1 \\ 0 \\ 0 \end{bmatrix} \quad [\hat{x}] = \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & -1 \\ 0 & 1 & 0 \end{bmatrix} = \begin{bmatrix} 1 & 0 & 0 \\ 0 & \cos\theta & -\sin\theta \\ 0 & \sin\theta & \cos\theta \end{bmatrix}$$

- Rotation matrix corresponding to $(1, 0, 1)^T \leftarrow \text{exp coordinate}$

$$\hat{\omega} = \begin{bmatrix} 1 \\ 0 \\ 1 \end{bmatrix} \cdot \frac{1}{\sqrt{2}}, \quad \theta = \sqrt{2}$$

Logarithm of Rotations

- If $R = I$, then $\theta = 0$ and $\hat{\omega}$ is undefined.

- If $\text{tr}(R) = -1$, then $\theta = \pi$ and set $\hat{\omega}$ equal to one of the following

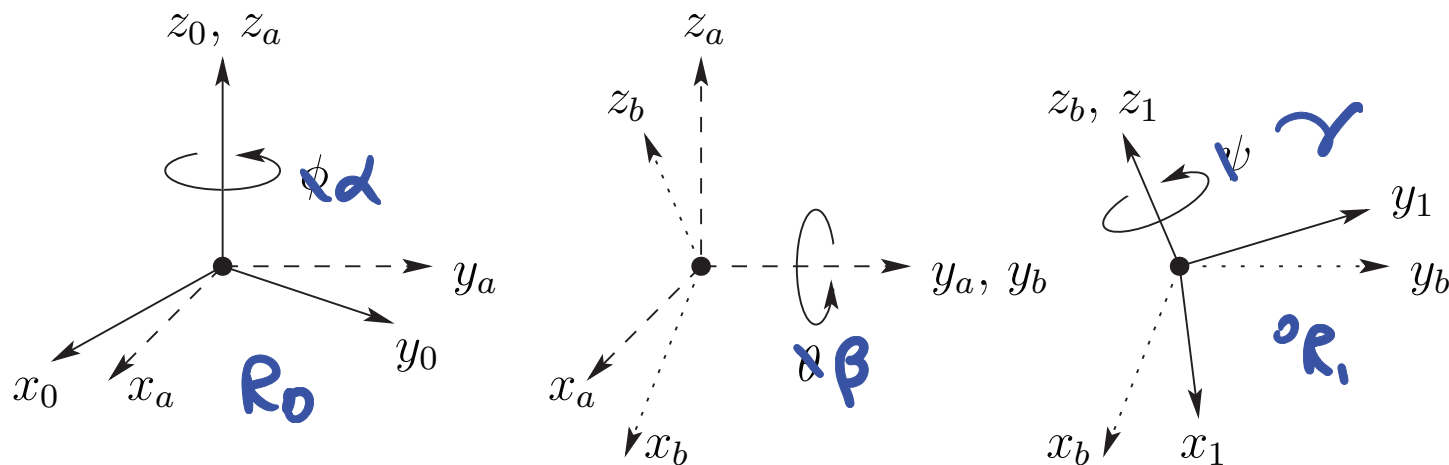
$$\frac{1}{\sqrt{2(1+r_{33})}} \begin{bmatrix} r_{13} \\ r_{23} \\ 1+r_{33} \end{bmatrix}, \frac{1}{\sqrt{2(1+r_{22})}} \begin{bmatrix} r_{12} \\ 1+r_{22} \\ r_{32} \end{bmatrix}, \frac{1}{\sqrt{2(1+r_{11})}} \begin{bmatrix} 1+r_{11} \\ r_{21} \\ r_{31} \end{bmatrix}$$

- Otherwise, $\theta = \cos^{-1} \left(\frac{1}{2}(\text{tr}(R) - 1) \right) \in [0, \pi)$ and $[\hat{\omega}] = \frac{1}{2 \sin(\theta)} (R - R^T)$

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Euler Angle Representation of Rotation



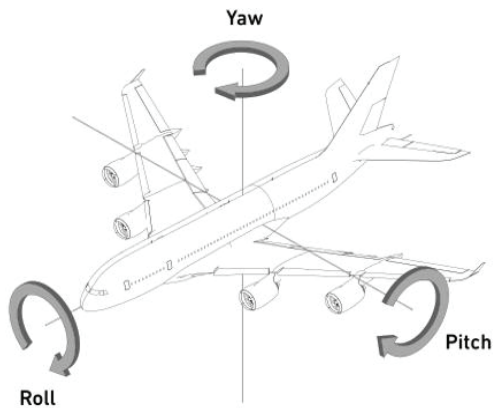
- A common method of specifying a rotation matrix is through three independent quantities called **Euler Angles**.
- Euler angle representation
 - Initially, frame $\{0\}$ coincides with frame $\{1\}$
 - Rotate $\{1\}$ about \hat{z}_0 by an angle α , then rotate about \hat{y}_a axis by β , and then rotate about the \hat{z}_b axis by γ . This yields a net orientation ${}^0R_1(\alpha, \beta, \gamma)$ parameterized by the ZYZ angles (α, β, γ)

$$- {}^0R_1(\alpha, \beta, \gamma) = \underline{R_z(\alpha)R_y(\beta)R_z(\gamma)} \quad \leftarrow \quad {}^0R_1(\alpha, \beta, \gamma) = e^{[\hat{\omega}]\theta}$$

$$\hat{\omega}\theta = \log({}^0R_1)$$

Other Euler-Like Parameterizations

- Other types of Euler angle parameterization can be devised using different ordered sets of rotation axes
- Common choices include:
 - ZYX Euler angles: also called *Fick angles* or yaw, pitch and roll angles
 $R_z(\alpha) R_y(\beta) R_x(\gamma)$
 - YZX Euler angles (Helmholtz angles)



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Exponential Map of $se(3)$: From Twist to Rigid Motion

Theorem 1.

For any $\mathcal{V} = (\omega, v)$ and $\theta \in \mathbb{R}$, we have $e^{[\mathcal{V}]\theta} \in SE(3)$

- Case 1 ($\omega = 0$): $e^{[\mathcal{V}]\theta} = \begin{bmatrix} I & v\theta \\ 0 & 1 \end{bmatrix}$
- Case 2 ($\omega \neq 0$): without loss of generality assume $\|\omega\| = 1$. Then

$$e^{[\mathcal{V}]\theta} = \begin{bmatrix} e^{[\omega]\theta} & G(\theta)v \\ 0 & 1 \end{bmatrix}, \text{ with } G(\theta) = I\theta + (1 - \cos(\theta))[\omega] + (\theta - \sin(\theta))[\omega]^2 \quad (1)$$

$$\mathcal{V} \in \mathbb{R}^{6 \times 1} = \begin{bmatrix} \omega \\ v \end{bmatrix}, \quad [\mathcal{V}] = \begin{bmatrix} \overset{3 \times 3}{[\omega]} & \overset{3 \times 1}{v} \\ 0 & 1 \end{bmatrix} \in \mathbb{R}^{4 \times 4}, \quad \underbrace{e^{[\mathcal{V}]\theta}} \in SE(3)$$

Forward map
exponential from $se(3)$ $\xrightarrow{\text{exp}(\cdot)}$ $SE(3)$
twist screw axis transformation matrix

Log of $SE(3)$: from Rigid-Body Motion to Twist

Theorem 2.

$$T = \begin{bmatrix} R & p \\ 0 & 1 \end{bmatrix}$$

Given any $T = (R, p) \in SE(3)$, one can always find twist $\mathcal{V} = (\underline{\omega}, \underline{v})$ and a scalar θ such that

$$e^{[\mathcal{V}]\theta} = T = \begin{bmatrix} R & p \\ 0 & 1 \end{bmatrix}$$

Matrix Logarithm Algorithm:

- If $R = I$, then set $\omega = 0$, $v = p/\|p\|$, and $\theta = \|p\|$.
- Otherwise, use matrix logarithm on $SO(3)$ to determine ω and θ from R . Then v is calculated as $v = G^{-1}(\theta)p$, where

$$G^{-1}(\theta) = \frac{1}{\theta}I - \frac{1}{2}[\omega] + \left(\frac{1}{\theta} - \frac{1}{2}\cos\frac{\theta}{2}\right)[\omega]^2$$

Exponential Coordinates of Rigid Transformation

- To sum up, screw axis $\mathcal{S} = (\omega, v)$ can be expressed as a normalized twist; its matrix representation is

$$[\mathcal{S}] = \begin{bmatrix} [\omega] & v \\ 0 & 0 \end{bmatrix} \in se(3)$$

- A point started at $p(0)$ at time zero, travel along screw axis \mathcal{S} at unit speed for time t will end up at $\tilde{p}(t) = e^{[\mathcal{S}]t}\tilde{p}(0)$
- Given \mathcal{S} we can use Theorem 1 to compute $e^{[\mathcal{S}]t} \in SE(3)$;
- Given $T \in SE(3)$, we can use Theorem 2 to find $\mathcal{S} = (\omega, v)$ and θ such that $e^{[\mathcal{S}]\theta} = T$.

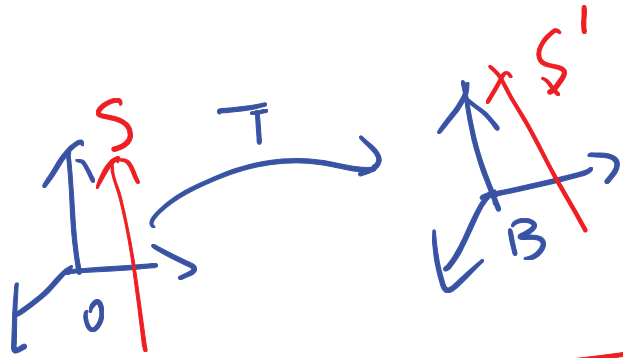


- We call $\mathcal{S}\theta$ the **Exponential Coordinate** of the homogeneous transformation $T \in SE(3)$

for any $T \in SE(3)$, $\Rightarrow [Ad_T] \cong \begin{bmatrix} R & 0 \\ [p]R & R \end{bmatrix}$
 $\hookrightarrow (R, p)$ 6×6

$$\underline{T = e^{[\hat{S}]\theta}} \rightarrow S \rightarrow S'$$

twist



• screw axis : - $S = (\hat{S}, h, q) \iff \begin{bmatrix} \omega \\ v \end{bmatrix} = \begin{bmatrix} \hat{S} \\ h\hat{S} - \hat{S} \times q \end{bmatrix}$
 $\dot{\theta} = 1$

- all rigid body motion can be "thought of" as screw motion
 rotation & linear motion along the same axis

- We typically write $V = S\dot{\theta}$

• Exponential coordinate & rotation operator (with $\{0\}$)

- ODE for rotation: $\dot{p} = \omega \times p = [\omega] p \Rightarrow p(t) = e^{[\omega]t} p(0)$

- if $\omega = \hat{w}$, unit vector, $t = \theta$

- $\hat{w}\theta \leftrightarrow R = e^{[\hat{w}]\theta} \in SO(3)$. $\hat{w}\theta$ is exp. coordinate of R

$$- p' = \underbrace{e^{[\hat{w}]\theta}}_R p$$

- Given frame $\{A\}$, $R_A = [\hat{x}_A, \hat{y}_A, \hat{z}_A]$, then $\underbrace{R R_A}_{\text{means rotate}}$ $\{A\}$ about \hat{w} by θ

- Expression of rotation operator R in another frame $\{B\}$.

Rotation matrix
in $\{0\}$

Rotation in $\{B\}$

- Exp coordinate for rigid body motion {with $\{0\}$ }

- ODE: $\dot{p} = v + \underbrace{[\omega \times]}_{[w]} p$ under twist $\mathcal{V} = \begin{bmatrix} \omega \\ v \end{bmatrix}$

- $\Rightarrow \dot{\tilde{p}} = \underbrace{\begin{bmatrix} [w] & v \\ 0 & 0 \end{bmatrix}}_{[V]} \tilde{p} \Rightarrow \tilde{p}(t) = e^{[V]t} \tilde{p}(0)$

- consider unit twist $[V] \in \mathbb{R}^{4 \times 4}$

- $SE(3) \xrightleftharpoons[\log]{\exp} SE(3)$
 $S \quad T = e^{[S]\theta}$

- $\tilde{p}' = T \tilde{p}$: rotate p along S by θ degree

- $T \cdot T_A$: rotate frame $\{A\}$ along ...

- rigid operation on screw axis: $[Ad_T] \cdot \underline{S}_1$ means rotate screw axis S_1 along S by θ degree

$T = e^{[S]\theta}$ \leftarrow $\underbrace{\quad}_{6 \times 6}$ $\underbrace{\quad}_{6 \times 1}$

Outline - expression of T in $\{B\}$

$$T \longleftrightarrow T_B^{-1} T T_B$$

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Instantaneous Velocity of Rotating Frame

- $\{A\}$ frame is rotating with orientation $R_A(t)$ and velocity $\omega_A(t)$ at time t
(Note: everything is wrt $\{O\}$ -frame)
- Let $\hat{\omega}\theta = \log(R_A(t))$ be its exp. coordinate.
 - Note: $\hat{\omega}\theta$ means $R_A(t)$ can be obtained from the reference frame (say $\{O\}$ -frame) by rotating about $\hat{\omega}$ by θ degree.
 - $\hat{\omega}\theta$ only describes the current orientation of $\{A\}$ relative to $\{O\}$, it does not contain info about how the frame is rotating at time t .
- What is the relation between $\omega_A(t)$ and $R_A(t)$?

$$\frac{d}{dt}R_A(t) = [\omega_A(t)]R_A(t) \Rightarrow [\omega_A(t)] = \dot{R}_A(t)R_A^{-1}(t)$$

$$R_A(t) = \begin{bmatrix} x_A(t) & y_A(t) & z_A(t) \end{bmatrix}$$

$$\dot{x}_A(t) = \omega_A \times x_A(t) = [\omega_A] x_A(t), \quad \dot{y}_A(t) = [\omega_A] y_A(t), \quad \dot{z}_A(t) = [\omega_A] z_A(t)$$

$$\Rightarrow \dot{R}_A = [\omega_A] R_A \Rightarrow [\omega_A] = \dot{R}_A(t) \cdot R_A(t)^{-1}$$

Instantaneous Velocity of Moving Frame

- $\{A\}$ moving frame with configuration $T_A(t)$ at time t undergoes a rigid body motion with velocity $\mathcal{V}_A(t) = (\omega, v)$ (Note: everything is wrt $\{O\}$ -frame)
- The exponential coordinate $\hat{S}\theta = \log(T_A(t))$ only indicates the current configuration of $\{A\}$, and does not tell us about how the frame is moving at time t .
- What is the relation between $\mathcal{V}_A(t)$ and $T_A(t)$?

$$\frac{d}{dt}T_A(t) = [\mathcal{V}_A(t)]T_A(t) \Rightarrow [\mathcal{V}_A(t)] = \dot{T}_A(t)T_A^{-1}(t)$$

$T_A = \begin{bmatrix} R_A & p_A \\ 0 & 1 \end{bmatrix}$, a frame can be determined by direction vector of axes, and origin p free vector point

In homogeneous coordinate: $\tilde{x}_A = \begin{bmatrix} x_A \\ 0 \end{bmatrix}$, $\tilde{y}_A = \begin{bmatrix} y_A \\ 0 \end{bmatrix}$, $\tilde{z}_A = \begin{bmatrix} z_A \\ 0 \end{bmatrix}$

$$\dot{\tilde{x}}_A = [\mathcal{V}_A] \tilde{x}_A, \quad \dot{\tilde{y}}_A = [\mathcal{V}_A] \tilde{y}_A \Rightarrow \dot{T}_A = [\mathcal{V}_A] T_A \Rightarrow [\mathcal{V}_A] = \dot{T}_A T_A^{-1}$$

More Space

More Space